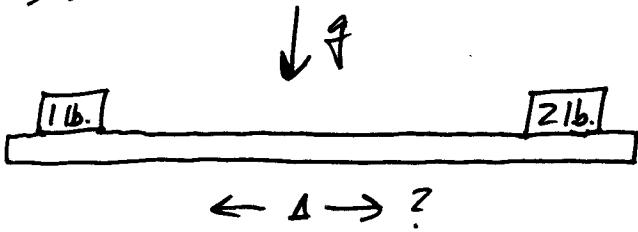


## Moments

In your physics class the physical quantity of moment of a force was most likely called a torque. They are the same thing.

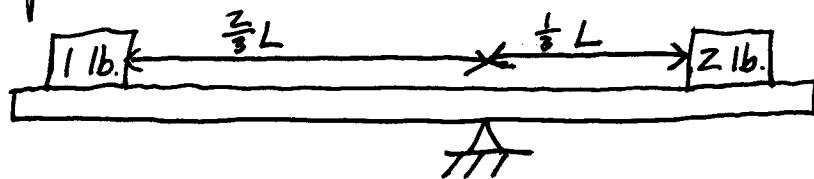
When we talk about a moment we need to specify a force, the point of application of that force, and the point which we are considering the moment of the force about.

Consider the following lever system (a teeter-totter).



for equilibrium

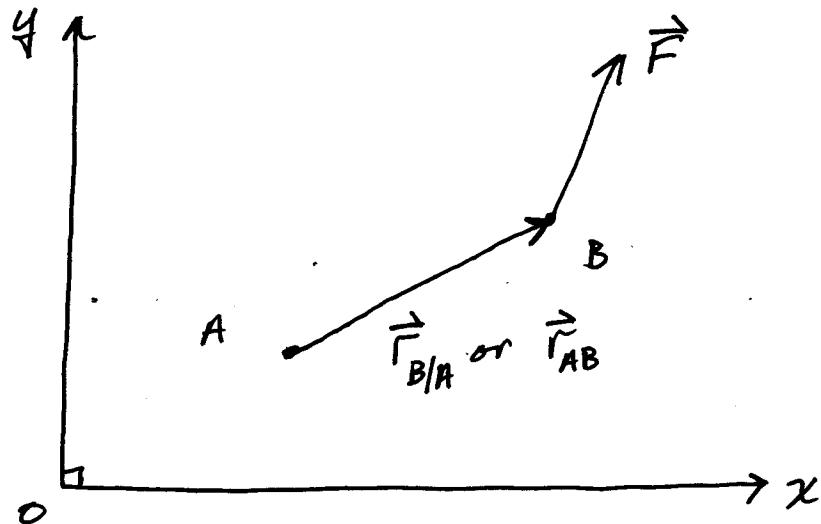
Where do we place the fulcrum? For those who have seen the balance of moments/torques before, we know that the fulcrum has to be placed closer to the 2 lb. weight such that the ratio of the distances to the fulcrum is equal to the ratio of the weights.



The moment of a force about a point gives the propensity of the force to cause a rotation about the point. The dimensions of the moment of a force are force multiplied by distance.

Looking at the lever, the 2 lb. weight tends to rotate the bar clockwise, and its moment about the fulcrum is  $2 \text{ lb.} \times \frac{1}{3}L$ . The 1 lb. weight tends to rotate the bar counterclockwise, and its moment about the fulcrum is  $1 \text{ lb.} \times \frac{2}{3}L$ . So, ultimately, the moments of these two forces about the fulcrum are equal in magnitude but tend to cause rotation in the opposite directions.

Let's consider a more general 2-D scenario.

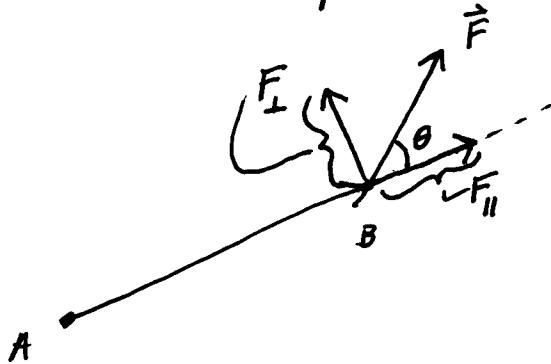


$\vec{r}_{B/A}$  = "position of B with respect to A"  
 $\vec{r}_{AB}$  = "position vector from A to B"

The question is, what is the moment of the force  $\vec{F}$ , which is applied at point B, about point A.

In general there is no need to have any structure connecting points A and B, but sometimes considering a bar (or a right-angled frame) connecting the two points is useful.

Let's break up the force  $\vec{F}$  into components perpendicular and parallel to  $\vec{r}_{AB}$ .



Does the  $F_{\parallel}$  component tend to rotate a bar connecting A and B? No, this component of the force tends to stretch  $\vec{r}_{AB}$ .

Then, the perpendicular component does tend to rotate such a bar and the magnitude of the moment is

$$M_A = r_{AB} \cdot F_{\perp}$$

Next, notice that  $F_{\perp} = F \sin \theta$ , so

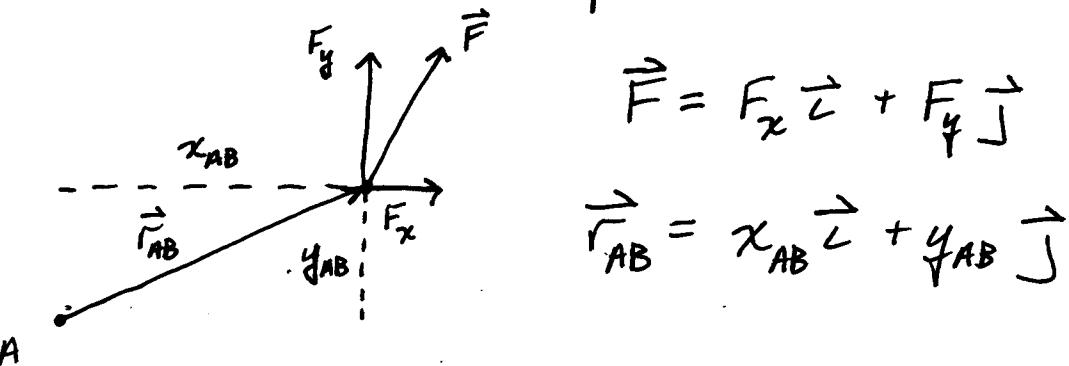
$$M_A = r_{AB} F \sin \theta$$

Finally, the direction of  $M$  is counterclockwise and if we use the right-hand rule this direction is perpendicular to both  $\vec{r}_{AB}$  and  $\vec{F}$  and is out of the page. So in vector form we have

$$\vec{M}_A = r_{AB} F \sin \theta \hat{e}_{\perp}$$

This looks like the cross product  $\vec{M} = \vec{r}_{AB} \times \vec{F}$ , but we need to analyze this in 3-D before we formally make this definition.

Let's look at this in components.



For the sign of the moment, using the right hand rule, we will call the direction out of the page the positive  $z$ -direction.

First consider the moment about point A due to the  $x$ -component of  $\vec{F}$ .

$F_x$  tends to "stretch" the  $x_{AB}$  component of the vector  $\vec{r}_{AB}$  but "rotates" the  $y_{AB}$  component in the clockwise ( $-z$ ) direction.

Similarly,  $F_y$  "stretches"  $y_{AB}$  and rotates  $x_{AB}$  counterclockwise ( $+z$ ).

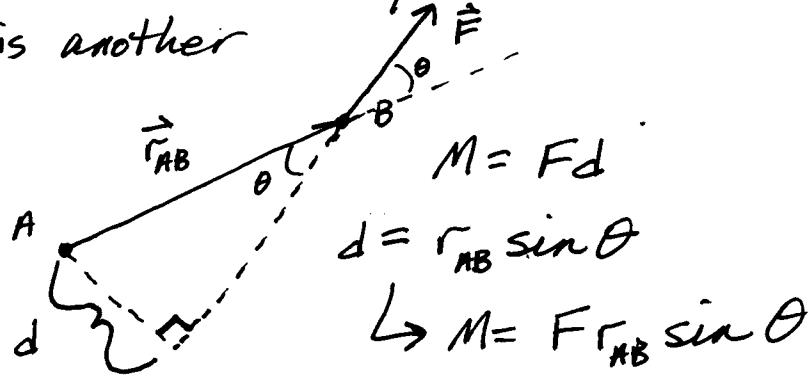
The total moment due to  $\vec{F}$  is the sum of these two contributions.

$$M = -F_x y_{AB} + F_y x_{AB}$$

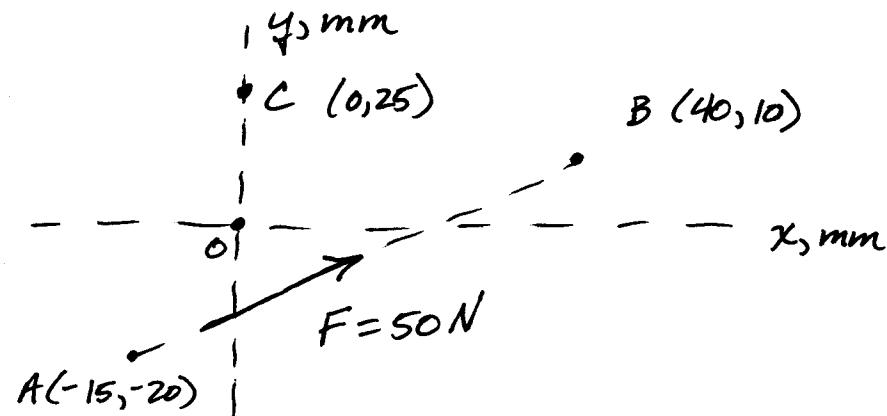
Terminology:  $y_{AB}$  is called the "moment arm" of the  $F_x$  component of the force about point A.

$x_{AB}$  is then the moment arm of the  $F_y$  component.

There can be several different ways to visualize moments in 2-D. Here is another

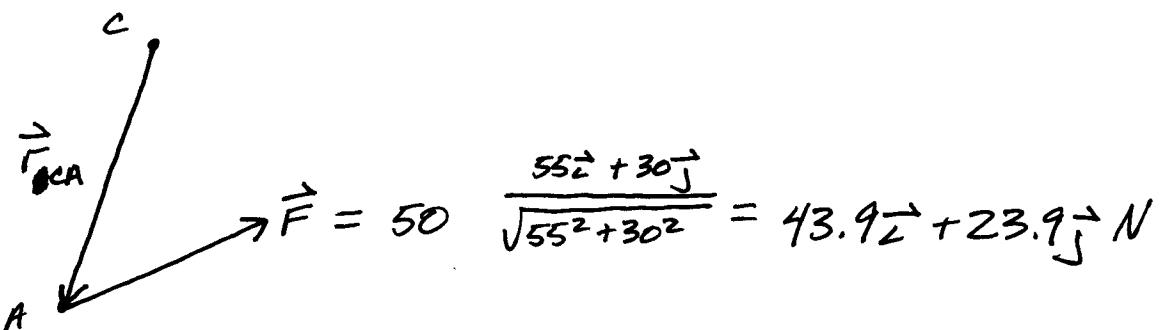


Example:



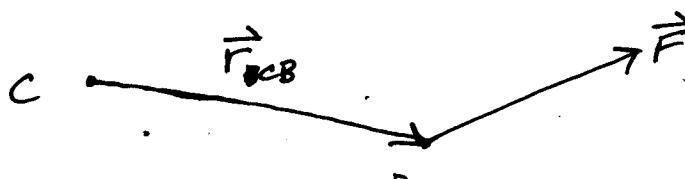
Determine the moment of the  $50\text{ N}$  force (a) about the point C as if it acts at A and (b) as if it acts at B.

a)



$$M_z^C = 43.9(45) - 23.9(15) = 1617 \text{ N}\cdot\text{mm}$$

b)



$$\vec{r}_{CB} = 40\hat{i} - 15\hat{j}$$

$$M_z^C = 23.9(40) + 43.9(15) = 1615 \text{ N}\cdot\text{mm}$$

Difference is due to truncation error.

(27)

We will ultimately show that in general vector form, the moment due to a force applied at point B about point A is

$$\vec{M}^{(A)} = \vec{r}_{AB} \times \vec{F}^{(B)} = \vec{r}_{B/A} \times \vec{F}^{(B)}$$

We can use this to prove the equality of our previous results.

$$(a) \quad \vec{M}^{(C)} = \vec{r}_{CA} \times \vec{F}^{(A)}$$

$$(b) \quad \vec{M}^{(C)} = \vec{r}_{CB} \times \vec{F}^{(B)}$$

Now recognize that the vector form of both  $\vec{F}^{(A)}$  and  $\vec{F}^{(B)}$  are the same, and let's just call this  $\vec{F}$ , i.e.

$$\vec{F}^{(A)} = \vec{F}^{(B)} = \vec{F}$$

$$\text{Next, let's write } \vec{r}_{CB} = \vec{r}_{CA} + \vec{r}_{AB}$$

$$(b) \rightarrow \vec{M}^{(C)} = (\vec{r}_{CA} + \vec{r}_{AB}) \times \vec{F}$$

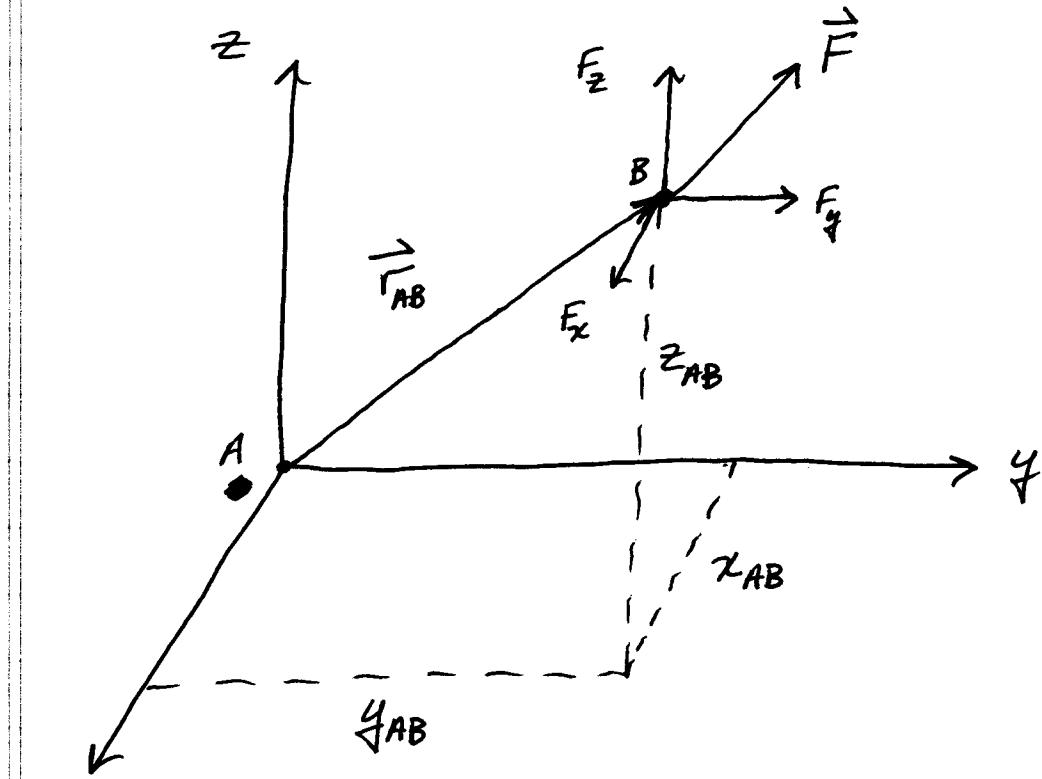
$$= \vec{r}_{CA} \times \vec{F} + \vec{r}_{AB} \times \vec{F}$$

$$\text{but } \vec{F} = F \frac{\vec{r}_{AB}}{|\vec{r}_{AB}|} \text{ and } \vec{r}_{AB} \times \vec{F} = \frac{F}{|\vec{r}_{AB}|} \underbrace{(\vec{r}_{AB} \times \vec{r}_{AB})}_{=0}$$

$$\rightarrow \vec{M}^{(C)} = \cancel{\vec{r}_{CA} \times \vec{F}} = \text{result from (a)}$$

(28)

## Moments in 3-D (translate origin to A)



$$\vec{F} = F_x \vec{i} + F_y \vec{j} + F_z \vec{k}$$

$$\vec{r}_{AB} = x_{AB} \vec{i} + y_{AB} \vec{j} + z_{AB} \vec{k}$$

The  $F_x$  component of the force rotates the  $z_{AB}$  component of the moment arm in the positive  $y$ -direction, and the  $y_{AB}$  component of the moment arm in the negative  $z$ -direction. This component of the force only tends to "stretch"  $x_{AB}$ .

$$\rightarrow F_x : F_x z_{AB} \vec{j} - F_x y_{AB} \vec{k}$$

$$\text{Similarly, } F_y : F_y x_{AB} \vec{k} - F_y z_{AB} \vec{l}$$

$$F_z : F_z y_{AB} \vec{l} - F_z x_{AB} \vec{j}$$

The sum of all of these contributions is the total moment of the force  $\vec{F}$  about point A.

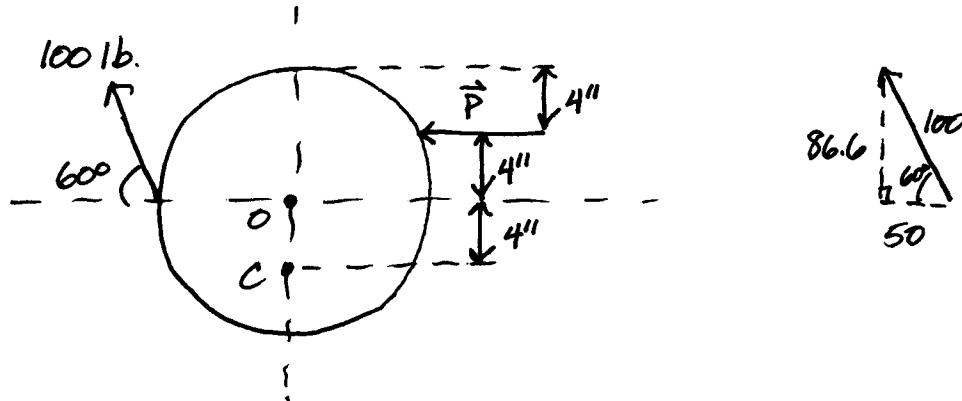
$$\begin{aligned} \vec{M}^{(A)} &= (F_z y_{AB} - F_y z_{AB}) \vec{l} \\ &\quad + (F_x z_{AB} - F_z x_{AB}) \vec{j} \\ &\quad + (F_y x_{AB} \vec{k} - F_x y_{AB}) \vec{k} \end{aligned}$$

Now consider  $\vec{r}_{AB} \times \vec{F}$

$$\begin{aligned} \vec{r}_{AB} \times \vec{F} &= \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ x_{AB} & y_{AB} & z_{AB} \\ F_x & F_y & F_z \end{vmatrix} = F_z y_{AB} \vec{l} + F_x z_{AB} \vec{j} + F_y x_{AB} \vec{k} \\ &\quad - F_y z_{AB} \vec{l} - F_z x_{AB} \vec{j} - F_x y_{AB} \vec{k} \\ &= \vec{M}^{(A)} \end{aligned}$$

$$\therefore \boxed{\vec{M}^{(A)} = \vec{r}_{AB} \times \vec{F}^{(B)}}$$

Back to a 2-D Example:



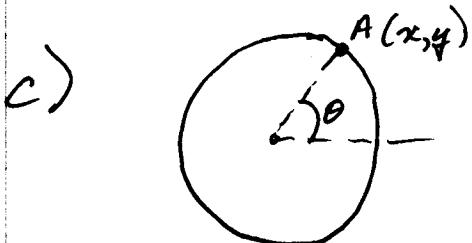
If the combined moment of the two forces about point C is zero, determine

- the magnitude of  $\vec{P}$
- the magnitude of  $\vec{R}$  the resultant of the two forces
- the  $x$  &  $y$  coordinates of a point A on the rim of the wheel about which the combined moment of the two forces is a maximum
- the combined moment  $M_A^C$  of the two forces about A

a)  $M_C^C = 0 = P \cdot 8 + 50 \cdot 4 - 86.6 \cdot 8 \rightarrow P = 61.6 \text{ lb.}$

b)  $\vec{R} = -50\vec{i} + 86.6\vec{j} - 61.6\vec{i} = -111.6\vec{i} + 86.6\vec{j}$

$$\rightarrow R = \sqrt{111.6^2 + 86.6^2} = 141.3 \text{ lb.} = R$$



$$x = 8 \cos \theta$$

$$y = 8 \sin \theta$$

$M_A^A = P(4-y) - 50y - 86.6(8+x)$ ,  $P=61.6$

$$\rightarrow M_z^A = -446.4 - 111.6y - 86.6x$$

$$= -446.4 - 111.6(8\sin\theta) - 86.6(8\cos\theta)$$

$$M_z^A = -446.4 - 892.8\sin\theta - 692.8\cos\theta$$

Maximize  $M_z^A$  with respect to  $\theta$

$$\frac{dM_z^A}{d\theta} = -892.8\cos\theta + 692.8\sin\theta = 0$$

$$\frac{\sin\theta}{\cos\theta} = \frac{892.8}{692.8}$$

$$\rightarrow \theta = 52.2^\circ, 232.2^\circ$$

$$M_z^A (\theta = 52.2^\circ) = -1577 \text{ lb-in} \quad (\text{min value but max magnitude})$$

$$M_z^A (\theta = 232.2^\circ) = \cancel{1577} 684 \text{ lb-in} \quad (\text{max value but } \underline{\text{not}} \text{ max magnitude})$$

Note: There will be a point on the rim where the moment is between these two extreme values with magnitude equal to zero.

Max magnitude at  $x = 8\cos 52.2^\circ = 4.9 \text{ in}$   
 $y = 8\sin 52.2^\circ = 6.3 \text{ in}$

d) Max magnitude is 1577 lb-in in the negative z-direction (clockwise).

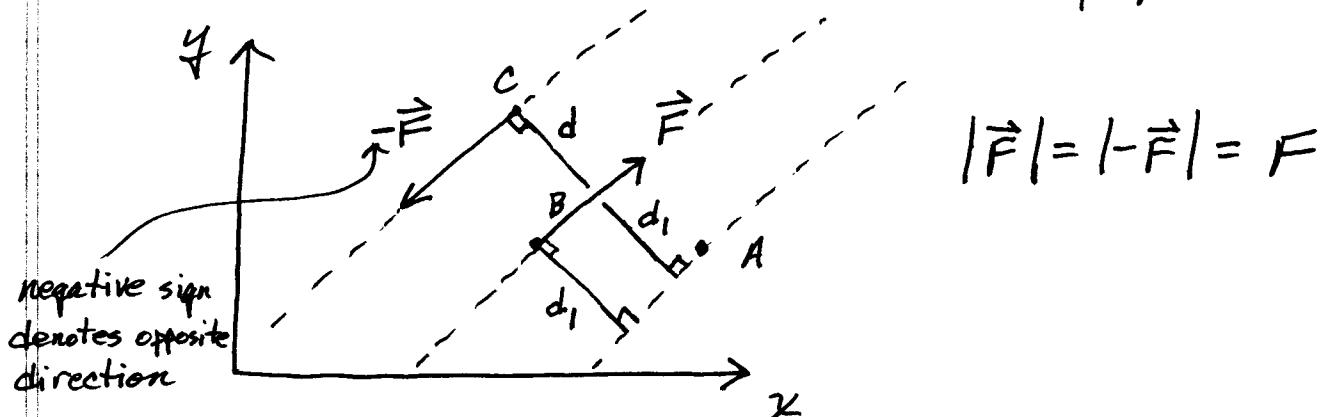
## Couples

A couple is a set of two equal and opposite non-colinear forces.

in magnitude

in direction

2-D : Consider the total moment due to a couple about an arbitrary point.



$$\begin{aligned}\sum M_A^A &= -Fd_1 + F(d_1 + d) \\ &= Fd\end{aligned}$$

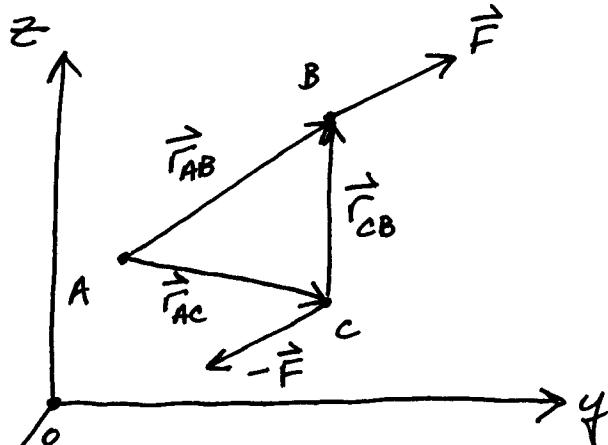
\* Note that this result has absolutely nothing to do with the location of point A.

\*\* Therefore, the moment due to a couple is the same about any point in space.

Equivalent representations of a couple:

$$F \swarrow d \searrow F = \curvearrowright C = Fd$$

3-D



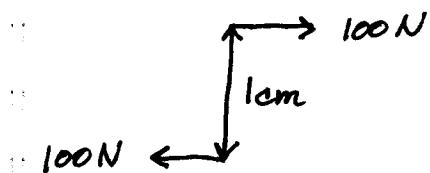
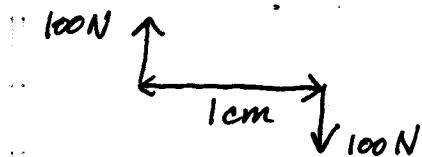
$$\text{Note: } \vec{r}_{AB} = \vec{r}_{AC} + \vec{r}_{CB}$$

$$\begin{aligned}\sum \vec{M}_A &= \vec{r}_{AB} \times \vec{F} + \vec{r}_{AC} \times (-\vec{F}) \\ &= (\vec{r}_{AB} - \vec{r}_{AC}) \times \vec{F}\end{aligned}$$

$$\vec{C} = \vec{r}_{CB} \times \vec{F}$$

$\vec{r}_{CB}$  has nothing to do with the location of point A.

### Equivalent Couples



$$\text{C } 1 \text{ N}\cdot\text{m}$$

## Equivalent Systems

Two sets of force systems are equivalent (as far as their effects on a rigid body are concerned) if the sum of all of the forces and the sum of all of the moments due to the forces and couples about any point in space are identical.

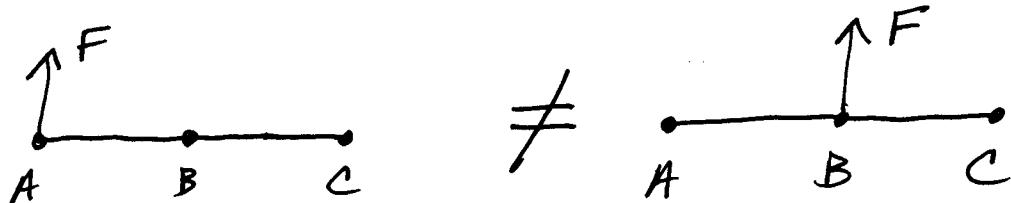
$$\sum \vec{F} = \sum_{i=1}^N \vec{F}_i$$

$$\sum \vec{M}^{(A)} = \sum_{i=1}^M \vec{C}_i + \sum_{i=1}^N \vec{r}_{AB_i} \times \vec{F}_i$$

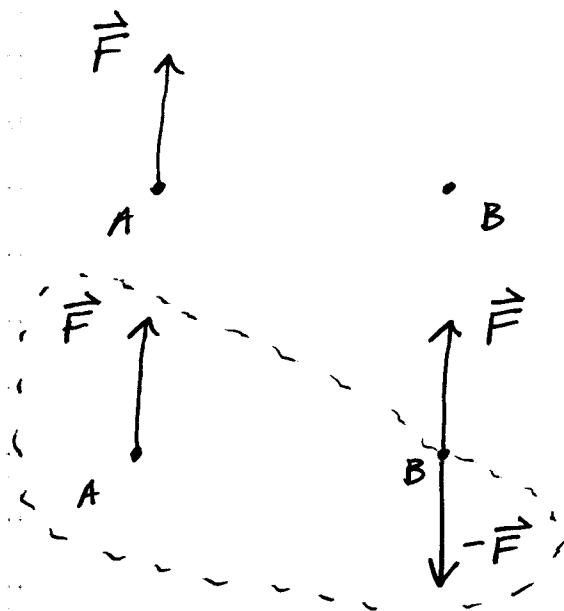
where there are  $N$  forces in the system acting at the points  $B_i$  and there are  $M$  arbitrarily positioned couples.

Recall that the net force of a couple is zero and the moment due to a couple about any point in space is the same. Therefore, a couple can be moved to any point in space without changing its effects on a rigid body.

However, changing the position of a force (not along its line of action) does change its moment about every point in space.



So, if we insist on moving a force, perhaps because it simplifies an analysis, we must compensate for the change in moment. Here is how this can be done:



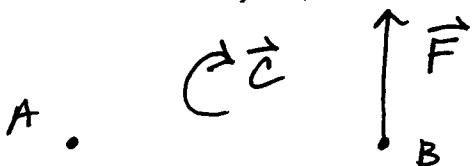
We want to move  $\vec{F}$  from A to B.

At B place both  $\vec{F}$  and  $-\vec{F}$ . Essentially we are adding zero to point B.

Now, consider the two forces contained within the dashed curve. This is a couple that mathematically is represented as

$$\vec{C} = \vec{r}_{BA} \times \vec{F}$$

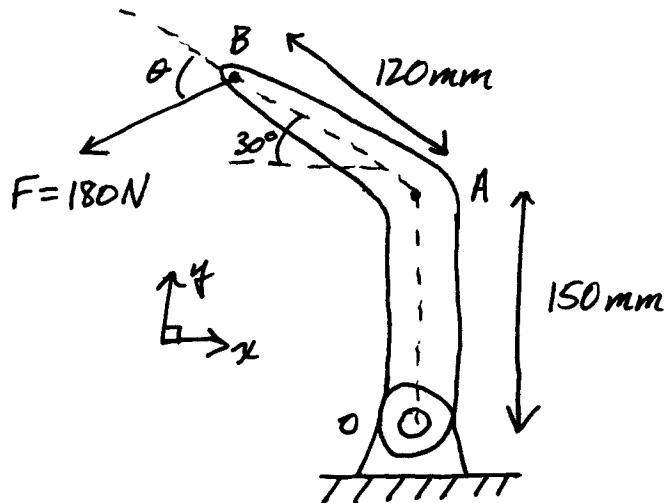
So the following is also an equivalent system



$\vec{F}$  acting at point B and a compensating couple  
 $\vec{C} = \vec{r}_{BA} \times \vec{F}$  acting at any point in space.

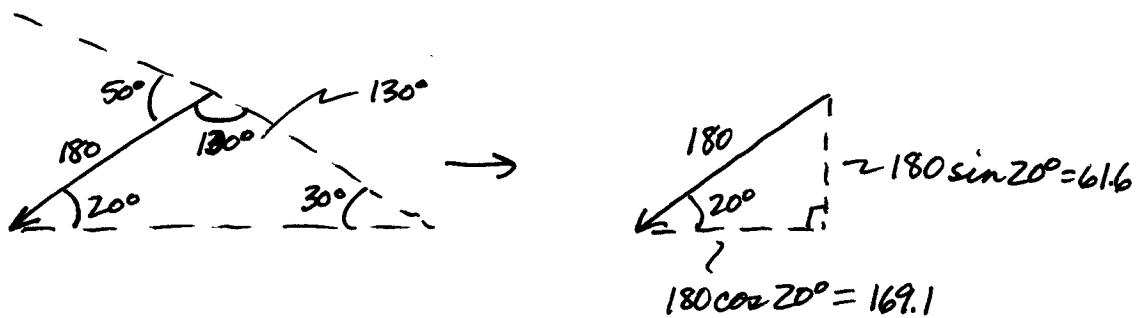
36

Example:



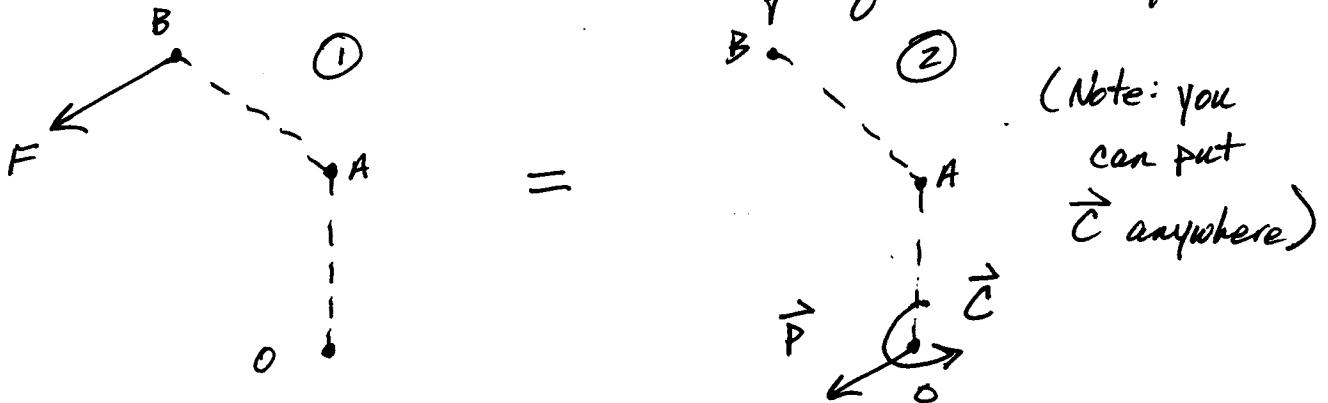
The 180-N force is applied to the end of the body OAB. If  $\theta=50^\circ$ , determine the equivalent force-couple system at the shaft axis O.

First, let's find the components of  $\vec{F}$ .



$$\rightarrow \vec{F} = -169.1 \vec{i} - 61.6 \vec{j}$$

We want to find the following equivalent system

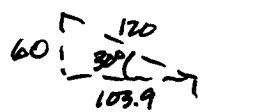


First, the sum of the forces in each system has to be the same,

$$\rightarrow \vec{F} = \vec{P} \rightarrow \boxed{\vec{P} = -169.1\vec{i} - 61.6\vec{j}} \text{ N}$$

Next, the moments about any point in space must be equivalent. Since the moment due to  $\vec{P}$  about O is zero it is easiest to choose this point.

$$\vec{r}_{OB} \times \vec{F} = \vec{C}$$

Vector method:   $\vec{r}_{OB} = (-103.9\vec{i} + 210\vec{j}) \text{ mm}$

$$\begin{aligned}\vec{C} &= (-103.9\vec{i} + 210\vec{j}) \times (-169.1\vec{i} - 61.6\vec{j}) \\ &= -103.9(-61.6) \underbrace{\vec{i} \times \vec{j}}_k + 210(-169.1) \underbrace{\vec{j} \times \vec{i}}_{-k} \\ &= 41,911 \vec{k} \text{ N} \cdot \text{mm} \text{ or } 41.9 \vec{k} \text{ N} \cdot \text{m}\end{aligned}$$

or  $C = 41.9 \text{ N} \cdot \text{m}$  counterclockwise (CCW)

Scalar method:  $F_x \text{ contribution} = 169.1(210) \text{ CCW}$   
 $F_y \text{ contribution} = 61.6(103.9) \text{ CCW}$

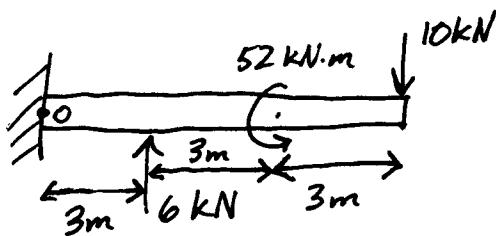
$$\rightarrow C = 169.1(210) + 61.6(103.9) = 41.9 \text{ N} \cdot \text{m CCW}$$

## Resultants

In my opinion finding the resultant of a set of forces is not a particularly useful procedure. However it is informative to know what the concept means. Definition: The resultant of a system of forces and couples is the simplest force - couple combination which can replace the original system without altering the effects on a rigid body. (Except for a distributed load.)

For 2-D systems the resultant is either a single force or a single couple.

Example: Determine and locate the resultant  $\vec{R}$  of the two forces and one couple acting on the beam.



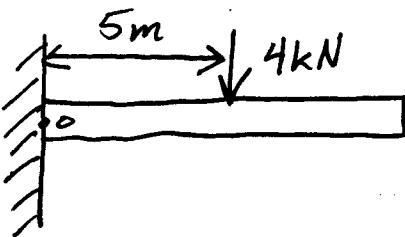
$$\sum F_x = 0$$

$$\sum F_y = -4 \text{ kN} \quad \text{or } 4 \text{ kN} \downarrow$$

$$\begin{aligned} \sum M_z^o &= 52 - 10 \cdot 9 + 6 \cdot 3 \\ &= -20 \text{ kN} \cdot \text{m} \end{aligned}$$

or  $20 \text{ kN} \cdot \text{m} \curvearrowleft$  (cw)

Resultant  $\curvearrowright$



$$\sum F_x = 0 \quad \checkmark$$

$$\sum F_y = -4 \quad \checkmark$$

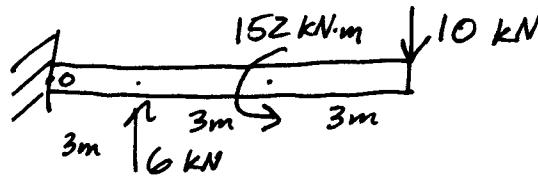
$$\sum M_z^o = -20 \quad \checkmark$$

What if we had:

$$\sum F_x = 0$$

$$\sum F_y = -4$$

$$\sum M_z^o = 152 - 10 \cdot 9 + 6 \cdot 3 = 80$$



Resultant  $\rightarrow$



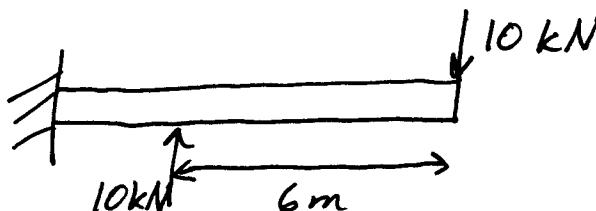
The resultant does not necessarily act at a point on the structure.

Last case:

$$\sum F_x = 0$$

$$\sum F_y = 0$$

$$\sum M_z^o = -60 \text{ kN} \cdot \text{m}$$



$\rightarrow$  Resultant is a couple  
 $\rightarrow$  acting anywhere in space

