



is split into two parts with each impinging on the surface of the specimen through a transparent viewing window, and then the positions of the reflected laser beams are tracked using two photo sensors. The distance between the two laser spots on photo sensors,  $S$ , is measured using the laser-positioning module. The bending curvature,  $\kappa$ , can be expressed as below [7]:

$$\kappa - \kappa_o = \frac{1}{R} - \frac{1}{R_o} = \frac{S_o - S}{2DL} \quad (1)$$

where  $R$  is the radius of curvature,  $D$  is the distance between the two laser spots on the specimen, and  $L$  is the distance between the specimen and photo sensors. The subscript “o” stands for the baselines of a reference state. The spatial relationships among these parameters are illustrated in Figure 2.

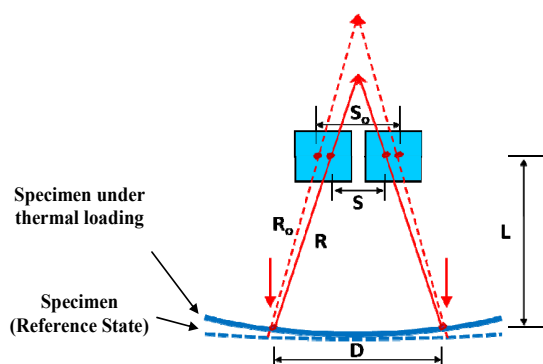


Figure 2. Spatial relationships among parameters in Equation 1.

### Bending beam specimen preparation: annular Cu via array in Si

Si wafers containing periodic arrays of blind annular Cu TSVs were co-designed by UT and siXis Inc. for bending beam measurement. To fabricate the annular vias, a square grid array of circular areas (90  $\mu\text{m}$  in diameter) is first photo-patterned on (001) Si wafers, and then a Bosch Si etching process is performed to drill circular holes to form periodic blind vias on the Si wafer. After the Si etching process, a conformal coating of dielectric material ( $\sim 2 \mu\text{m}$  thick silicon oxide) is deposited on the wafer surface and at the via sidewall. Finally, the wafers are electroplated with a thin layer of Cu  $\sim 2 \mu\text{m}$  in thickness to form annular Cu vias.

The samples used for bending beam measurements are diced from the Si wafer into 4 mm-by-40 mm strips. In order to measure the thickness of Cu on the via sidewalls, the bending beam specimens are polished on the edge to expose the cross-section of a row of TSVs. A cross-section SEM image of the annular via array is shown in Figure 3.

The electroplated annular Cu layer does not cover the whole length of the via sidewalls. As shown in Figure 3, the electroplated Cu film covers about upper two thirds of the via sidewalls. The dimensions of annular vias are measured from the cross-section of the specimen using a SEM/FIB system. Table 1 lists the averaged height and thickness of the annular via structure.

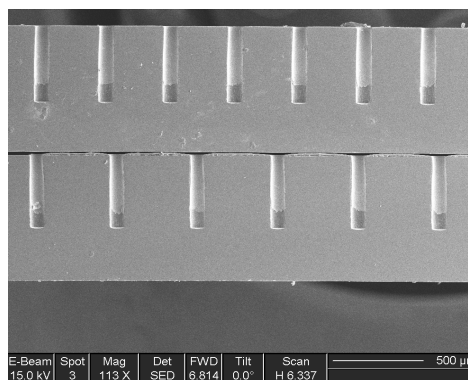


Figure 3. Cross-section image of bending beam specimen.

Table 1. Bending beam sample dimension.

Annular via array ( $D_f$ : 90 $\mu\text{m}$ , pitch: 360 $\mu\text{m}$ )	Unit: $\mu\text{m}$
Thickness of Wafer	$710 \pm 3$
Height of via hole	$429 \pm 5$
Height of Cu via	$311 \pm 7$
Thickness of annular Cu via	$1.6 \pm 0.3$
Thickness of silicon oxide	$2.0 \pm 0.2$

To complete the TSV test structure, the overburden layers of silicon oxide and Cu on the wafer surface are mechanically polished off. The top Si surface is exposed after polishing. The finished bending beam specimen is illustrated in Figure 4.

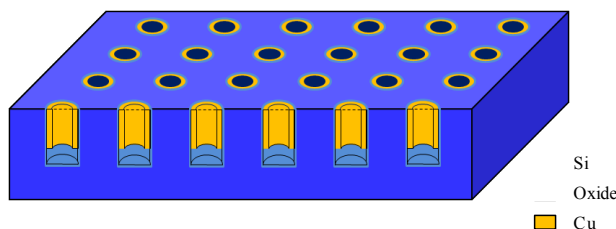


Figure 4. Visual illustration of bending beam specimen.

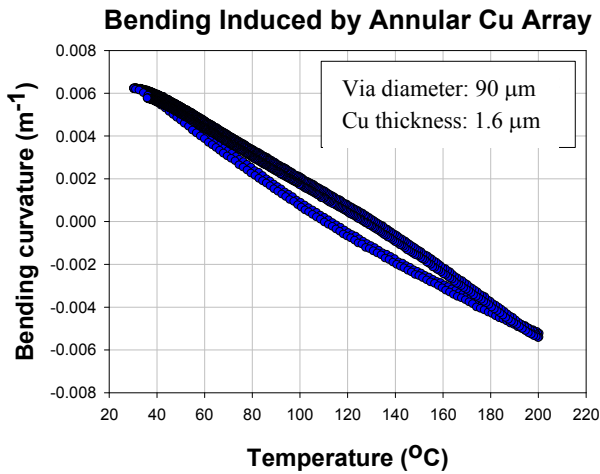
### Bending beam experiment and results

After the sample preparation, a specimen containing vias of 90  $\mu\text{m}$  in diameter is selected for the bending beam experiment. The pitch size between adjacent vias is 360  $\mu\text{m}$ . The specimen is placed in a vacuum test chamber with a 100 Torr of ambient dry nitrogen. The specimen is placed upside down on a supporting stage so that the incident laser beams are reflected by the backside of the specimen, i.e. the side without exposed vias. If the front side (i.e. the side containing via array) is selected as the reflecting surface, incident laser beams will be diffracted by the periodic via array, yielding multiple diffracted laser spots on the photo sensors and false signals in the position indicator.

To start the stress measurement, the bending beam specimen is thermal-cycled between the room temperature and 200°C in the vacuum chamber. The bending curvature of the specimen is continuously monitored as a function of temperature. Several factors contribute to the bending of the specimen: (a) the thermal expansion of Cu vias, (b) the thermal expansion of silicon oxide liner, and (c) the initial

curvature of the bare Si beam. To extract the thermal stress contribution of the Cu vias to the beam bending, the baseline curvatures induced by (b) and (c) have to be subtracted out from the overall bending. For this purpose, Cu vias are etched off from the test specimen using a nitric acid solution, and the baseline curvature is measured again under the same thermal cycles. By subtracting off the baseline curvature, the bending induced by the thermal expansion of the Cu via array is deduced.

Following this procedure, the bending curvature induced by Cu vias is obtained, and the result is plotted in Figure 5 as a function of temperature. The result shows clearly that the deformation of the Cu vias exhibits hysteresis phenomenon during two thermal cycles between 30°C and 200°C. Since Si and silicon oxide are both linear elastic in nature, the hysteresis behavior can be attributed to plasticity of the Cu layer in the via. This will be further discussed in the next section when the observed bending curvature in Figure 5 is analyzed to deduce the stress characteristics of the via structure.



**Figure 5.** Bending curvature induced by periodic annular Cu array under repeated thermal cycles between 30°C to 200°C.

### Thermal stress analysis for annular Cu TSVs

In this study, the thermal stress characteristics are deduced from the bending curvature of the Cu via structures observed during thermal cycling. First, an analogous example of bending beam stress measurement on a thin film is described as following. Consider an isotropic thin film lying on top of a thick Si substrate. The thermal stress on the thin film as a result of thermal loading is equal-biaxial and 2-D where there is no out-of-plane stress. Assuming that the thin film thickness  $h_f$  is much thinner than the substrate thickness  $h_s$ , the equi-biaxial stress of the thin film,  $\sigma_f$ , is proportional to the net curvature induced by the film,  $\kappa - \kappa_o$ :

$$\sigma_f = \frac{B_s h_s^2}{6h_f} (\kappa - \kappa_o) \quad (2)$$

where  $B_s$  is the biaxial modulus of the substrate. Equation 2 is the well-known Stoney's equation [8].

Now consider the case of annular Cu vias embedded in the Si matrix (Figure 4). Inside an annular Cu via, the thermal expansion mismatch between Cu and the Si matrix can induce thermal stresses during thermal cycling along the radial direction ( $\sigma_{rr}$ ), the circumferential direction ( $\sigma_{\theta\theta}$ ), and the axial direction ( $\sigma_{zz}$ ). The stress state in such a Cu via is 3-D in nature, and the bending curvature of the whole Si beam is induced by the radial stress  $\sigma_{rr}$  and the circumferential stress  $\sigma_{\theta\theta}$ . Due to the complex geometry, the stress components cannot be deduced simply from the bending curvature based on Equation 2. Therefore, an FEA is used to calculate the stress behavior of the Cu TSV test beam and to establish the relationship between the bending curvature  $\kappa$  and the stress components during thermal cycling. Based on the observed bending curvature, the stress components of the Cu via can be deduced. The FEA simulation procedure is described in the following section.

### FEA simulation for annular Cu TSVs

The thermal stress-induced bending curvature is simulated using a commercial software (ANSYS®). A simplified FEA model containing a 5-by-2 square grid annular Cu via array is constructed, as illustrated in Figure 6a. The geometrical dimensions of Cu vias in this model are listed in Table 1. Note that the height of annular Cu vias is less than half the thickness of Si wafer (Figure 3). A 3-D solid element (SOLID185) in the ANSYS software is selected for simulations. The elemental mesh in one of the ten unit cells is shown in Figure 6b, while the elemental mesh of the annular Cu/silicon oxide is magnified and shown in Figure 6c. The thermo-mechanical properties of materials used for simulations are listed in Table 2. Among them, the CTE of silicon oxide is intentionally assumed to be the same as that of Si. By excluding the CTE mismatch between silicon oxide and the Si substrate, the beam bending due to the thermal expansion of Cu material is simulated and is to be compared with the experimental results in the next section.

A symmetric boundary condition is applied on the surfaces at  $x = 0$  and at  $y = 0$ . The FEA model is simulated under varying thermal load  $\Delta T$  (-50°C ~ -200°C). After the simulation is complete, the bending curvature and the stress components are extracted from the result files for analysis.

**Table 2.** Thermo-mechanical properties of materials [9,10,11]

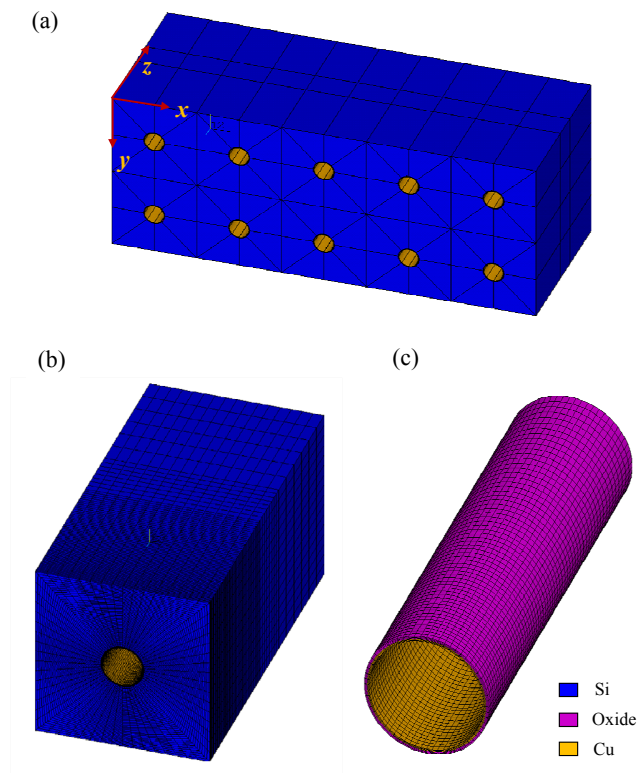
Material	CTE (ppm/K)	Young's Modulus (GPa)	Poisson's ratio
Si	2.3	130	0.28
Cu	17	110	0.35
Silicon oxide	2.3 <sup>1</sup>	76	0.18

### Thermal stresses on top Si surface

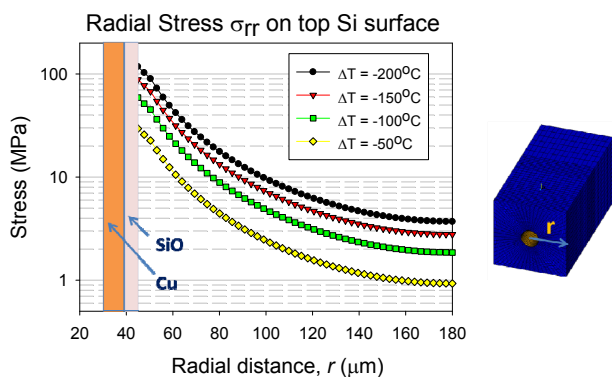
The thermal stresses on the top Si surface are first analyzed. In a cylindrical coordinate system, the stress components are axi-symmetric. According to the semi-analytical near-surface solution [2], the radial stress  $\sigma_{rr}$  and the circumferential stress  $\sigma_{\theta\theta}$  on the top Si surface (Figure 6a,  $z = 0$ ) are of the same magnitude but opposite in sign. The out-of-plane stress component  $\sigma_{zz}$  vanishes on the surface.

<sup>1</sup>The CTE of silicon oxide is intentionally assumed to be the same as that of Si.

Thus, only the radial stress from FEA simulation is examined here. The distribution of  $\sigma_{rr}$  on the top surface obtained by FEA is plotted as a function of the radial distance  $r$  in Figure 7. Under thermal load  $\Delta T$ , the bending curvature of the FEA model increases monotonically with the increasing thermal load. The radial stress distribution also increases uniformly with the bending curvature on the top Si surface.



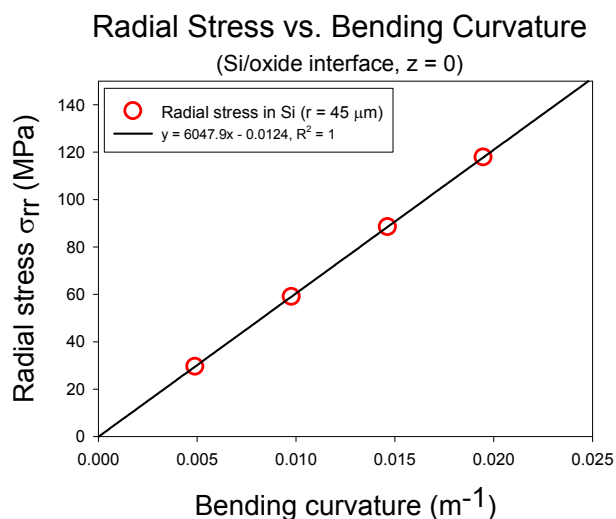
**Figure 6.** (a) FEA model containing a 5-by-2 annular Cu via array. (b) Element mesh in one unit cell. (c) Element mesh in annular Cu via and the silicon oxide liner.



**Figure 7.** Distribution of radial stress  $\sigma_{rr}$  on the top Si surface with varying thermal load  $\Delta T$ .

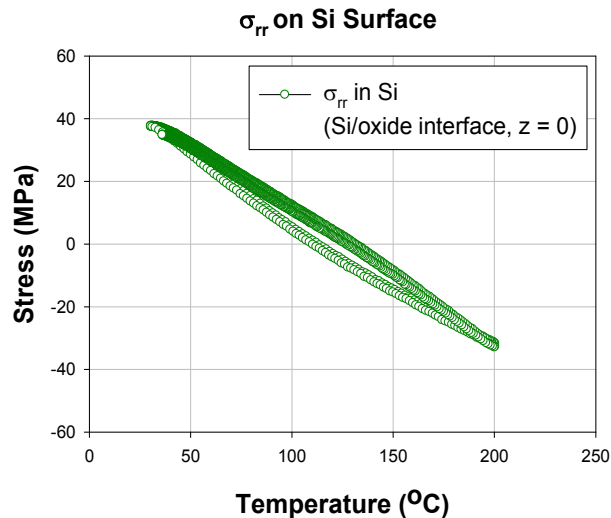
To show that the radial stresses in Figure 7 are proportional to the bending curvature, the simulation result at the TSV/Si boundary on the top surface ( $z = 0$ ) is selected and plotted against the curvature (Figure 8). It exhibits a linear relationship between the radial stress and the bending curvature. Therefore, the linear curve-fit equation shown in

Figure 8 can be used to obtain the corresponding radial stress once the bending curvature of the test beam is experimentally determined.



**Figure 8.** A linear relationship between the bending curvature and the radial stress  $\sigma_{rr}$  at a specific location (top Si surface, Si/oxide interface).

After establishing the stress-curvature relationship using the FEA, the bending curvature from the experiment is converted to the radial stress for the selected location ( $r = D_p/2, z = 0$ ). This corresponds to the circumference of the Cu TSV on the top Si surface where the radial stress can be deduced from the measured beam curvature (Figure 5) and the result is plotted in Figure 9 as a function of temperature for two repeated thermal cycles. As shown in Figure 9, the maximum radial stress  $\sigma_{rr}$  in Si is within a range of 39 MPa (at 30°C) to -33 MPa (at 200°C). The bending beam specimen has a stress-free temperature of 112°C or 125°C, depending on whether a positive or negative thermal load is applied. The non-linear behavior of stress-temperature curve results from the plasticity of Cu material during thermal cycles, which is discussed in the next section.

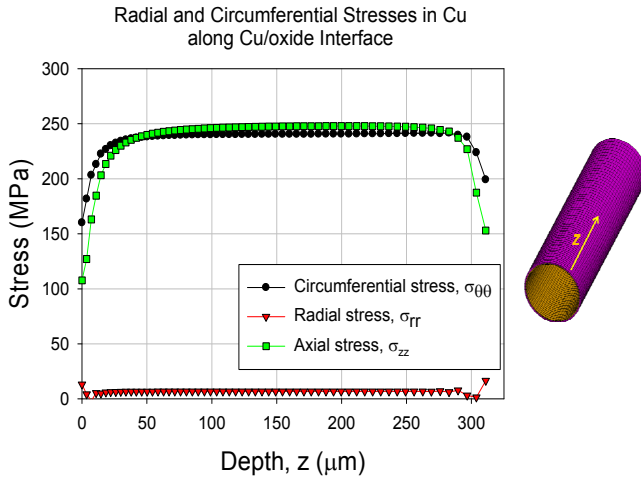


**Figure 9.** Radial stress  $\sigma_{rr}$  vs. temperature (top Si surface, Si/silicon oxide interface).

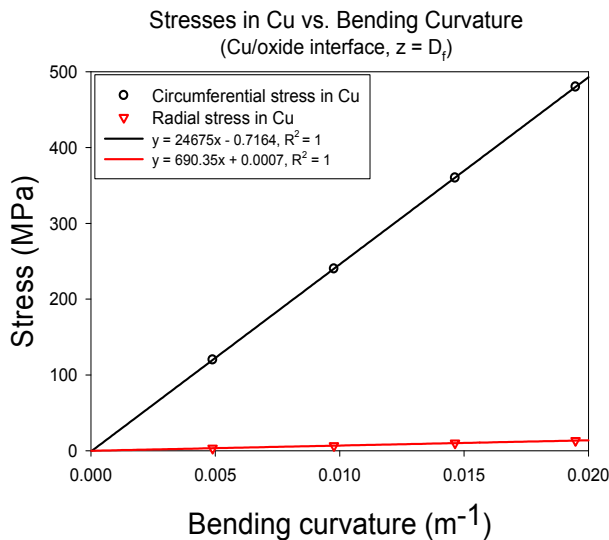


### Thermal stresses in annular Cu TSVs

Here the FEA-simulated thermal stresses in an annular Cu via are analyzed. In a cylindrical coordinate system, the thermal stress distribution in the annular Cu via is also axisymmetric. Under a thermal load  $\Delta T = -100^\circ\text{C}$ , thermal stresses in Cu along the Cu/oxide interface are plotted as a function of depth  $z$  in Figure 10.



**Figure 10.** Stresses in annular Cu as a function of depth  $z$  (Cu/oxide interface,  $D_f = 90 \mu\text{m}$ ).



**Figure 11.** Linear relationship between the bending curvature and the circumferential stress  $\sigma_{\theta\theta}$  and the radial stress  $\sigma_{rr}$  in annular Cu via (Cu/oxide interface).

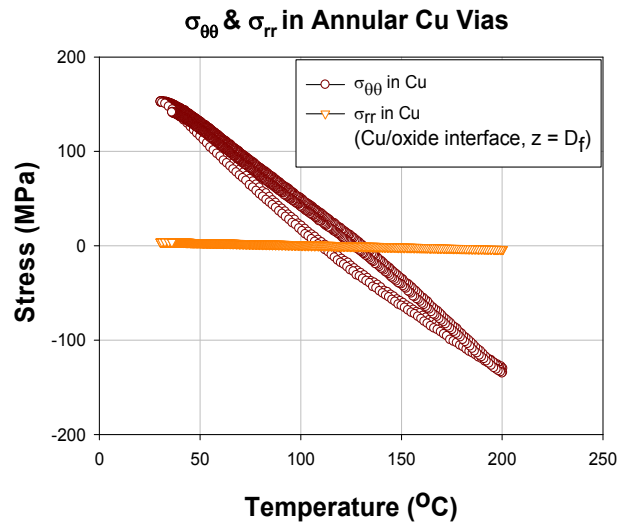
As shown in Figure 10, the circumferential stress  $\sigma_{\theta\theta}$  and the axial stress  $\sigma_{zz}$  are of a similar magnitude, but the radial stress  $\sigma_{rr}$  is smaller by a factor of 20. This can be attributed to the fact that the thickness of annular Cu via is much thinner than the via diameter  $D_f$  so that  $\sigma_{rr}$  is equivalent to the normal stress in a thin Cu film along the out-of-plane direction. It is interesting to note the stress variation as a function of the depth where the stresses approach a steady-state at a distance  $\sim D_f$  away from both ends of the Cu cylinder. Therefore, over most of the annular Cu via, except near the ends, the stress

state may be approximated by a 2-D plane-strain solution for a thin-wall cylinder.

In the following, the stress components which contribute to beam bending ( $\sigma_{\theta\theta}$  and  $\sigma_{rr}$ ) are chosen for analysis. FEA results at a depth  $z = D_f$  on the Cu/oxide interface are selected for this discussion. FEA-simulated thermal stresses in Cu are plotted as a function of the bending curvature in Figure 11. Linear equations between stresses and the bending curvature are obtained by curve fitting.

The stresses in annular Cu vias at  $z = D_f$  can be deduced based on the linear relations established by FEA simulations (Figure 11) and the bending curvature observed (Figure 5). It should be pointed out that such linear relationships between the stresses and the bending curvature applies no matter whether the Cu material has been mechanically yielded or not. The bending curvature of the Si beam originates from the force balance between the Si beam and the via structures on top. Because the Si material is linear elastic, the force (or stress) required to bend the Si beam is linearly proportional to the bending curvature.

The radial and circumferential stresses at the chosen location ( $r = D_f/2 - t_{\text{oxide}}$ ,  $z = D_f$ ) are plotted as a function of temperature in Figure 12.



**Figure 12.** Stress-temperature curves of annular Cu vias (Cu/oxide interface).

In Figure 12, the radial stress  $\sigma_{rr}$  in the annular Cu via is almost zero. It is expected because of the traction-free surface on the inner sidewall of the via. The circumferential stress  $\sigma_{\theta\theta}$ , which is much larger than the radial stress, ranges from +150 MPa (at  $30^\circ\text{C}$ ) to -130 MPa (at  $200^\circ\text{C}$ ). It implies that only the circumferential stress (or force) in Cu vias accounts for the specimen bending, because the other stress component  $\sigma_{rr}$  is a small quantity and  $\sigma_{zz}$ , although of a larger quantity, does not contribute to bending. The other observation from Figure 12 is the elastic hysteresis behavior of Cu. The annular Cu via is subjected to a thermal load followed by a reverse thermal load. During the reverse loading (i.e. cooling), the annular Cu via exhibits an initial elastic unloading behavior, followed by a so called “early yielding” occurring at a stress level below 50 MPa. Early yielding is often observed in thermal stress measurements for metallic thin films, which is similar to the

Bauschinger effect in bulk metals. Dislocation theories were applied in the past to explain such early yielding behavior of Cu thin films [12]. Because the annular Cu via investigated here is only ~1.6  $\mu\text{m}$  thick, the annular Cu via is equivalent to a Cu thin film. The observed “early yielding” behavior may be the same as that of Cu thin films.

### Summary

A stress measurement on Cu via arrays using bending beam technique has been demonstrated in this paper. The state-of-the-art curvature measurement technique has a sensitive resolution of  $10^{-3} \text{ m}^{-1}$  in curvature, making it possible to detect the beam deflection induced by a small magnitude of stress in Si ( $< 50 \text{ MPa}$ ). It has been found challenging to measure the stresses in such annular via structure using other techniques such as micro-Raman.

Bending beam technique can also be used to observe the plasticity behavior of Cu material during thermal cycling. In Figure 12, a non-linear elastic behavior has been observed in the electroplated Cu material, and the stress-free temperature of Cu vias can be determined from such stress-temperature curve. Compared to micro-Raman measurements, which are often operated at room temperature, the bending beam technique provides useful temperature-dependent information for thermo-mechanical analyses. However, the major disadvantage of the bending beam method is the sample configuration, which requires a large population of periodically arranged vias. The other disadvantage is that only averaged global stresses/strains of the whole beam are measured. Such a technique cannot measure stresses near individual vias as micro-Raman does. The local stresses in individual vias are subjected to vary with respect to Cu thickness, oxide thickness, the elastic anisotropy in Cu grains, etc. The precise local stresses induced by a specific TSV cannot be obtained using FEA simulations with only averaged sample geometry and with isotropic material properties. Instead, TSV-induced local stresses near the Si surface can be measured using micro-Raman. Consequently, for one to characterize the thermal stresses in TSV structures, the bending beam method and micro-Raman can be employed together as complementary techniques.

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